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Stimulated Emission Cross Section in Xenon-Neon Lasers

Gulalla Yaseen Bakr

Physics Department, College of Education, Salahaddin University, Salahaddin, Iraq.

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Abstract:

Gas Lasers are important tools that are used in variety purposes, for their low and (cw) output power. The aim of this study was to prepare a way to calculate an optimum stimulated emission cross-section in a gas laser containing a mixture of Xenon and Neon by (30%-70%). The process was a theoretical study of each gas in separate in terms of their physical properties as an active medium. The results of these calculations are logic and more convenient than other mixtures used before.

Key words: Gas-mixtures, Ionization and Excitation processes, Gas lasers

Introduction:

When gases are excited from their ground state several processes face this excitation, like collisions, scating, absorption, transfer and ionization[1]. Out of these, in gases the process of emission is the matter of question in gas lasers. Since the most considerable process is the stimulated emission which results the process of amplification and power coupling. A laser has a resonator which contains usually an active medium to be amplified then[2] the consequent processes mentioned above own its special cross section. The algebraic sum of these give what is called the total cross section[3]. The two main portions of cross section are excitation and the momentum transfer cross sections. Studies prove that these are the back bone of both spontaneous and stimulated emission cross sections.

Theory:

The excitation cross section results through elastic atom-atom collision with cross section $\sigma_e = \frac{4\pi}{k^2} (2l + 1) \sin^2 \delta_l(k)$ [4].

$$f(r, k) = \delta_l(k). k \dots (1)$$

Here, k is the usual wave vector $\left[\frac{2\pi}{\lambda} \right]$ (l is the impact factor [either zero or one] depending on the collision direction and $\delta_l(k)$ is a modules called the partial wave amplitude the M.T.C.S (Mixed Tabulated Complex Series. on the other hand is raised as the result of different impact angle and it is defined as the cross section produced through momentum exchange between the colliding species defined by[5]

$$\sigma_{MT} = \frac{4\pi}{k} \text{Im} f(k, \theta = 0) \dots (2)$$

Where Im is the imaginary part of the complex wave amplitude.

Analysis of equation (1) and (2) will be through lengthily [but straight forward

Bessel function of the first kind [6] and using Matlab2 program the results are

$$\sigma_{ex} = \frac{4\pi}{K^2} \left[\tan^{-1} \frac{(Ka)^{(2l+1)}}{(2l+1)(2l-1)!} \right] \dots (3)$$

Here (a) is a constant depending on the energy of excitation.

Finally the M.T.C.S will be manipulated and using the same category as for σ_{ex} , we obtain:

$$\sigma_{MT} = \frac{4\pi}{K^2} \left[1 - \frac{K^2}{(1-2l)} + \frac{\left(\frac{k^2}{2}\right)^2}{2!(1-2l)(3-2l)} - \frac{\left(\frac{k^2}{2}\right)^3}{3!(1-2l)(3-2l)(5-2l)} \dots \right] \dots (4)$$

The values of these cross sections are shown in Figures (2),(3) and (4) respectively.

Stimulated Emission cross section is defined through the analogy of absorption cross section on which is the rate of absorption of energy per atom[6].

$$\sigma_{ab} = \frac{4\pi^2 \alpha h^2}{m\nu} |Ref(ck)| \dots (5)$$

But the stimulated emission cross section is deferent in mechanism[7].

In gas lasers σ_{st} is formulated through the atomic transition characteristics.

In fact, it's the product of population inverse difference and the popular radiation width

$$\rho(\nu) = \left[\frac{c^2}{8\pi\nu^2} A_{12}(\nu) \right] \dots (6)$$

The 1st one is $[N_2 - N_1] \frac{g_2}{g_1}$ and A_{21} is called Einstein coefficient for spontaneous emission, from this definition[7].

$$\sigma = \frac{4\pi^2 \alpha h^2}{m\nu} \lambda_{21} A_{21} [N_2 - N_1] \frac{g_2}{g_1} \dots (7)$$

α is constant called the fine structure constant: $\frac{1}{(137)^2}$, λ is the wavelength of radiation \hbar is $\left(\frac{h}{2\pi}\right)$ the famous Planks constant, g_1 and g_2 are called the statistical weight of the upper and lower atomic state $(2J+1)$ were J is total angular momentum m is the mass of atomic contributing species (here is Xenon) since Neon is just activating material [8].

The energy level diagram of the mixture is shown in figure number(1).

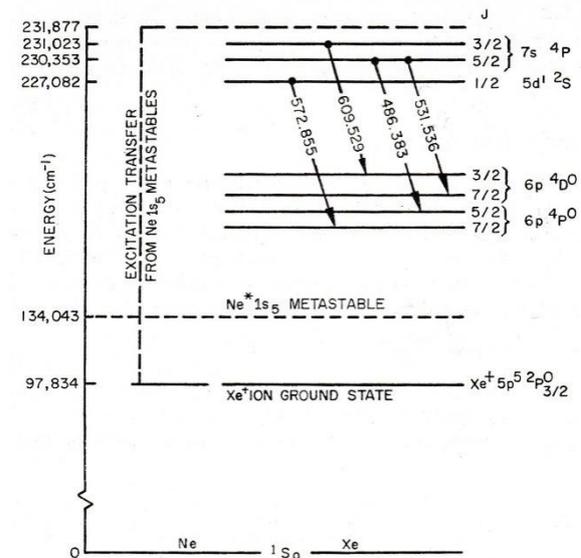


Fig.(1):Energy level diagram for Xenon and Neon gases[2]

Results and Discussion:

A mixture has been used with the following properties throughout the whole work, we tried to overcome all problems concerning different excitation processes facing a mixture of xenon and neon gases by the ratio of 30:70[Xe:Ne] in a discharge tube hoping that it will be used as an active medium for an ion laser known as Xenon-Neon ion laser.

The parameters are important to be considered and they are the ionization potentials of these gases as 12.13eV and 21.56eV as in Figure(1)respectively. These two values give the idea that Xe ionizes (9.43)eV before its counterpart Neon, so the resonant transfers for Xenon from Neon at this value gives the possibility of (Xe) to have almost many excited states after ionization is well known that two main processes[atomic and ionic] are in question:

a. The excitation cross –section for both gases in the range of (0-20)eV for Xe , and (0-50)eV for Ne respectively. The results of these calculations are shown in Figure(2).

b. Avery popular and famous counter part of the first one is the

momentum transfer cross –section with in the range mentioned earlier, the results of these variations are shown in Figure (3)for the mixture fixed as (30:70) percent for (Xe:Ne) gases.

B.Brandsen and Joachain[9] have proposed the ionization cross – section for both gases at their peak values to be $(0.113 \times 10^{-16})\text{cm}^2$ for Xenon and $(0.0333 \times 10^{-16})\text{cm}^2$ for neon respectively. Our values in Figure(5) give them a $(0.11 \times 10^{-16}$ and $0.032 \times 10^{-16})\text{cm}^2$ for the gases which are in good agreement with the experimental values described.

c. Stimulated emission cross-section found both theoretically and experimentally give what is called [saturation output intensity if they are used as an active medium for a medium powered Ion laser workable at (5371\AA) was:

$$I_{sat} = \frac{h\nu_L}{\tau_m \sigma_{st}}$$

Her ν_L is the frequency of the operating laser wave length. τ_m is the spontaneous life time at the 5p state and σ_{st} is the desired stimulated emission cross-section.

$$\text{Here } \nu_L = c/\lambda = \frac{3 \times 10^8}{0.5371 \times 10^{-6}} = 5.58 \times 10^{14} \text{Hz.}$$

$$\tau_s = 8.2 \text{nsec. } \sigma_{st} = 0.0333 \times 10^{-20} \text{m}^2.$$

As Such:

$$I_{sat} = \frac{6.625 \times 10^{-34} \times 5.58 \times 10^{14}}{8.2 \times 10^{-9} \times 0.0333 \times 10^{-20}} = 135.38 \frac{w}{m^2}.$$

The theory of Boltzmann transverse mixtures predicated that at these circumstances, Inert gases are to be considered as cold-plasma and Enhanced by Electrical methods.

d. Finally, much work have been done to find out the proposed stimulation emission cross – section for these gases in their peak positions since, in laser, this parameter plays an important role for the output power they were $(0.3 \times 10^{-16})\text{cm}^2$ and $(24 \times 10^{-16})\text{cm}^2$ for Ne and Xe respectively.

Now, recommended tool to be confident about the results is that we have to ascertain about the results of figure (4) which represents the sum of stimulated and ionization cross –section taking the values of pure ionization potential from this result, we obtained $(0.1333 \times 10^{-16})\text{cm}^2$ for the mixture which is quite reasonable with the value obtained by [9] to be $(0.1323 \times 10^{-16})\text{cm}^2$.

Conclusion:

In this study the tendency was towards a formulation for what is called stimulated emission cross-section in a mixture containing Xenon and Neon gases. The test was through fixing a ratio which is common in the field of cold plasma parameter calculations. The facts obtained in the process are:

The best fitting ratio for manipulating the mixture thermodynamically was (Xenon: Neon) as (30:70) as usual. This ratio is coming from the fact that at this ratio, the resonance transition between Ne 2p and Xe 5p are in its highest activity.

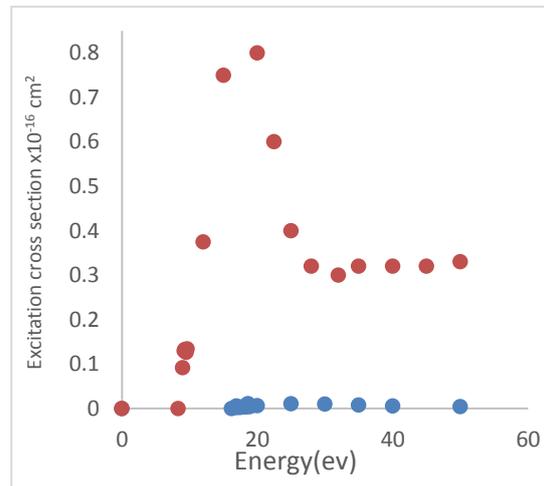


Fig.[2].Excitation cross section for Xenon and Neon as a function of energy

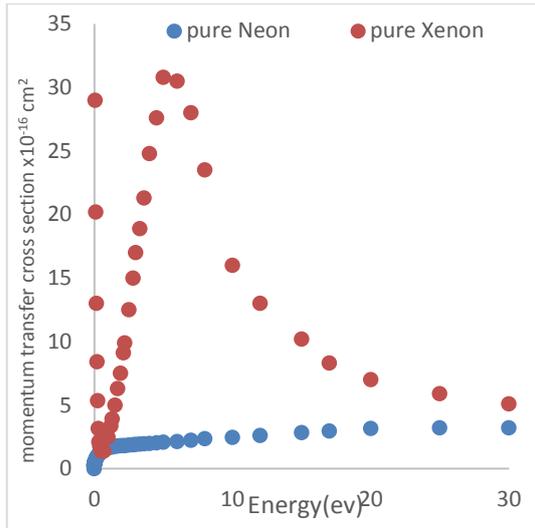


Fig.[3].Variation of momentum transfer cross section as a function of energy for each of Xe and Ne.

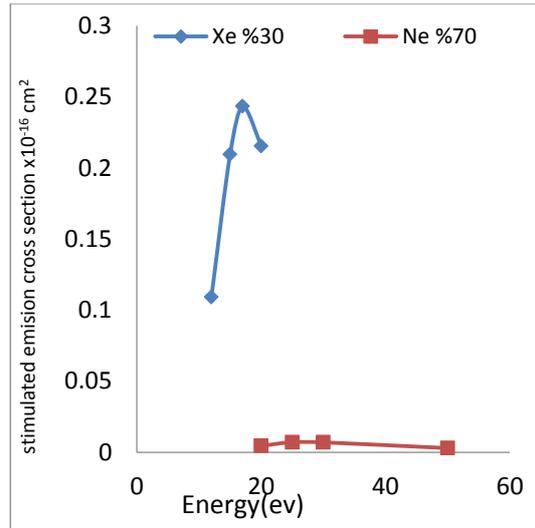
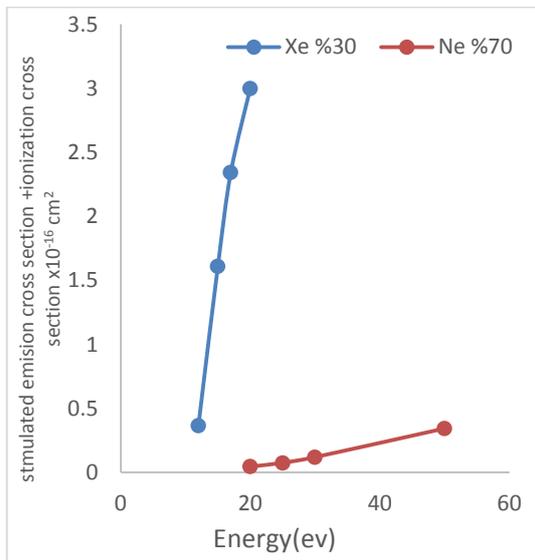


Fig.[5].stimulated emission cross section for Xenon Neon gas mixture as a function of energy.



Fig[4].combination of stimulated emission cross section +ionization cross section for xenon and neon gases as a function of energy.

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مساحة مقطع الانبعاث المحفز في الليزر نيون – زينون

غولاله ياسين بكر

قسم الفيزياء ، كلية التربية ، جامعة صلاح الدين، صلاح الدين، العراق.

الخلاصة:

تعتبر الليزرز الغازية من التقنيات المهمة التي تستخدم في عدة مجالات و سبب ذلك هو قدرتها الواطئه و ذات الموجه المستمرة. والهدف من هذه الدراسة هو حساب وتمائل ما يسمى مساحة مقطع الانبعاث المحفز في ليزر غازى يملك خليط من غازى النيون والزينون بالنسبة (30:70)% على التوالى. وتمت هذه من خلال دراسة نظرية لكل غاز على حدة من حيث الصفات الفيزيائية تعمل كوسط فعال وكانت النتائج منطقية واكثر مناسبة من الدراسات السابقة.

الكلمات المفتاحية: خليط الغازات، عمليات التاين والتهيج، الليزرز الغازية.