

Calculating the Radii of the Even-Even Nickel Isotopes

Safa Mohsin Ibrahim ^{1*}, Abdallah Ben Rhaiem ¹, Sameera Ahmed Ebrahiem ³

¹Department of Physics, Faculty of Sciences, University of Sfax, Sfax, Tunisia.

²Department of Physics, College of Education for Pure Sciences (Ibn Al-Haitham), University of Baghdad, Baghdad, Iraq.

*Corresponding Author.

Received 17/02/2024, Revised 17/05/2024, Accepted 19/05/2024, Published Online First 20/07/2024



© 2022 The Author(s). Published by College of Science for Women, University of Baghdad.

This is an open-access article distributed under the terms of the [Creative Commons Attribution 4.0 International License](https://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

Abstract

The current investigation concentrated on even-even nuclei forms to isotopes that have masses number less than 100 ($A > 100$) to ($^{58-68}_{28}Ni$) elements. This includes a look at deformation parameters (β_2) derived from reduced electric transition probability ($E2$) \uparrow as well as distortions parameters (δ) derived by intrinsic electric quadrupole moment (Q_0). A Roots mean square Radius $\langle r^2 \rangle^{1/2}$. We studied the most important deformation parameters (δ, β_2), and major and minor of ellipsoid axes (a,b) in addition to the difference between these axes (ΔR). All these parameters were calculated using the deformed shell model equations and applied theoretically in a special program. Differences in nuclei forms were observed for selected isotopes via drawing two 2-dimensional shapes. According to the current findings. Thus the distortion coefficients decrease as the number of neutrons approaches the magic number. Thus it was observed that the most distorted and inclined isotopes of the elliptical shape are the isotopes with numbers of protons and neutrons far from the magic numbers. Also, the obtained results were compared with the theoretical results from Raman source by "single-shell Asymptotic Nilsson Model" (SSANM), and noticed that was little change in the results. Also observed from the results that ($^{58}_{28}Ni$) is the most stable isotope and has a clear spherical shape because (Z) is equal to (28) a magic number and (N) is equal to (30) a number close to the magic number.

Keywords: A deformation parameters (β_2, δ), An electric quadrupole moment Q_0 , A probability of the transition $B(E_2: 0^+ \rightarrow 2^+)$, A Roots mean square Radius $\langle r^2 \rangle^{1/2}$.

Introduction

The protons and neutrons that make up the shell structure of the atomic nucleus are mirrored by the shape of the nucleus ¹. If the shells are filled, we are talking about a "magic" nucleus, which is spherical in shape ². However, most cores tend to be that deformed because their shells are only partially filled. The most common shapes encountered are oblate (flat, pancake and pillow) or prolate (elongated, cigar- and rugby-ball-shaped) ³. In some

cases, rearranging the neutrons or protons within the nucleus itself is enough to change a shape, so a nucleus can take various shapes depending on different energies ⁴. If the energy between these states is close, the other shapes can combine ⁵. Depending on the relative pivot values of the ellipsoid shapes, it can subdivide deformed nuclei into oblate, prolate, and triaxial deformed nuclei. It can be observed that there are different shapes of

nuclei in the ground state and only a few have a spherical shape, predominating in the same nucleus 5. The arrangement of valence nucleons in an empty shell is what causes the nuclear deformation, hence a deformation, only happens whenever both the proton

(p) & neutron (n) shells are half filled ⁶. Quadruple deformations parameter which is degree of nucleus shape difference from the sphere, as shown in Fig. 1 ⁷.

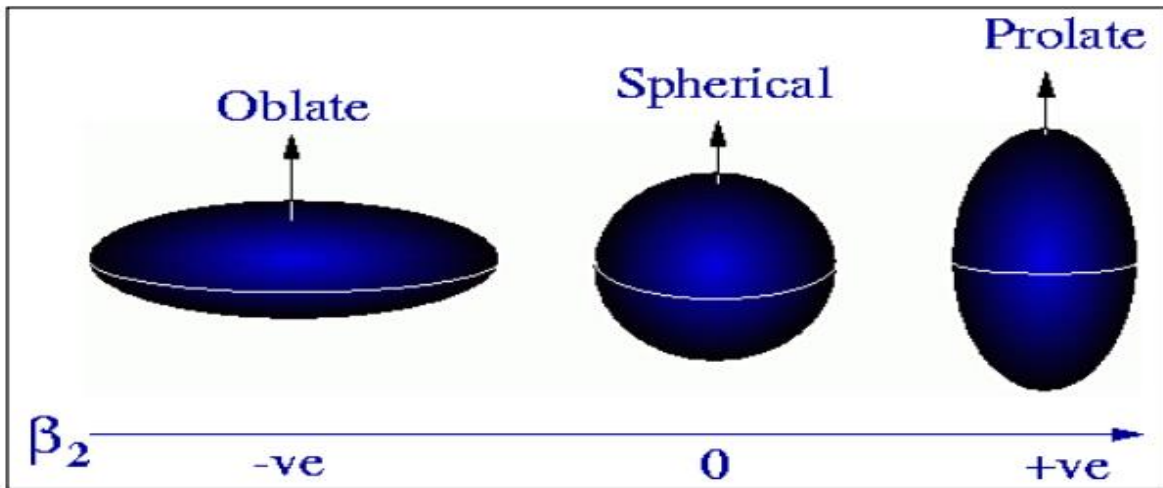


Figure .1. Diagram showing oblate, spherical and prolate shapes. The arrows for the oblate and prolate shapes indicate the symmetry ⁷.

The main objective of the study is to calculate the radii of even-even nickel isotopes by using the special deformation equation and applying it theoretically in matlab programs, and through it also calculate the deformation coefficients, through which the shape of the nuclei was determined.

Deformation parameter:

According to the liquid-drop model, nuclei may be pliable and soft, this indicates that the shape of the nucleus may differ significantly from the spherical shape. The researchers found a large number of nuclei in regions with neutron number (N) and proton number (P) away from the magic values that have a noticeably distorted charge distribution ⁷.

The assumption of independent nucleon motion in an average potential well serves as the theoretical underpinning of the chance model. the specific nucleon's interactions with every, other nucleon in nucleus are represented by this potential well ⁸.

The transition probability B(E2) is used to calculate the deformation parameter β_2 for the $2^+ \rightarrow 0^+$ transition. In the second (2^+) and first excited states 0^+ , there is an alternative to quadratic analysis. Basic nuclear information that supplements our

understanding of nuclide level energies ^{8,9}. The (E2) \uparrow values, in general, Confirmation of the large range, of quaternary, distortion in nuclides, estimated using the Global Beast Fit equation:

$$B(E2) \uparrow = 2.6 E\gamma^{-1} Z^2 A^{-2/3} \dots\dots\dots 1$$

, Where: Z is the atomic number,

A is the mass number of a nucleus, and

$E\gamma_0$ energy of gamma ray transitions in units of KeV.

This indicates, that (E2) values in that equation are fundamental on experimentally-measured amounts ($E\gamma, Z, A$) instead of relying on nuclear models ¹⁰. The parameter of the deformation β_2 is an effective, model that depends on quantum mechanics and is simple to visualize¹⁰. Assuming this potential well describes the interactions of a particular nucleon with all other nucleons in the nucleus ¹⁰:

$$\beta_2 = ,4\pi / 3ZR_o^2 [B(E2) \uparrow e^2 b^2 / e^2]^{1/2} \dots\dots\dots 2$$

Where: Ro: is the nuclear average radius, which is derived from the equation:

$$R_o^2 = 0.0144 A^{2/3} \text{ barn} \dots\dots\dots 3$$

The deformation parameter δ is calculated using intrinsic quadrilateral electrical moments Q_0 , which is used as it provides important information about shape and deformation because it measures the deviation of the charge distribution from spherical symmetry¹¹.

$$\delta = 0.75Q_0 / (\langle r^2 \rangle > Z) \dots\dots\dots 4$$

, Through the following equations, the mean radius $\langle r^2 \rangle$: can be calculated

$$\langle r^2 \rangle = 0.63R_0^2 (1 + 10/3(\pi a_0/R_0)^2) / (1 + (\pi a_0/R_0)^2) \quad (A \leq 100) \dots\dots\dots 5$$

$$\langle r^2 \rangle = 0.63(1.2A^{1/3})^2 \quad (A > 100) \dots\dots\dots 6$$

Where: R_0 : The radii potential parameters are $R_0 = 1.07 A^{1/3}$ fm and, $a_0 = 0.55$ (fm), and a_0 is derived from the information on fast electron scattering. An self-quadruple moment, Equation-based calculation via⁹:

$$Q_0 = [(16 \pi / 5) B(E_2) e^2 b^2 / e^2]^{1/2} \dots\dots\dots 7$$

Nuclear Shape

When nuclei are stable their shape is generally spherical. The goal of this effort reduces the energy of surface. As a result, little spherical sections are noticed such as in the region $150 < A < 190$ and only the ratio can be utilized to define these irregularities¹²:

$$\delta = \Delta R / R \dots\dots\dots 8$$

, Where:

ΔR denotes the variance between the semi-, minor and semi-, major axis

R is an average of the nuclear radius.

$$\Delta R = (b - a) \dots\dots\dots 9$$

Where:

a, b : major and minor of ellipsoid axes.

radius of the root mean square charge (isotope transformation):

(rms) is the radius of the nuclear charge $R = \langle r^2 \rangle^{(1/2)}$, together with other nuclear ground state

properties, and is considered one of the keys to information about the main nuclear materials that indicate the effectiveness of the nuclear structure, for example: closure crust and start to deform¹³.

, Direct inference of the root square radius (rms), $\langle r^2 \rangle^{(1/2)}$, is made from the scattering electron distribution; The radius of charge distribution for an evenly charged sphere is squared^{14,15}:

$$\langle r^2 \rangle = \frac{3}{5} R^2 = \frac{3}{5} R_0^2 A^{2/3} > 100 \dots\dots\dots 10$$

Where:

A : number mass

R : spheres radius.

$$R = R_0 A^{1/3}$$

A probability of reduced electric quadrupole transition (E_2) ↑

The best way to create a nuclear structure and to test models of nuclear structures is through radiated electromagnetic transitions between nuclear states¹⁶. The magnitude of the critical electric quadruple moments the energy of the lower nucleation levels the average periods of nucleation states and the nuclear distortion coefficient β depend on $B(E_2)$, transition. The collective effects, that multiple, nucleons can be involved in are referred to as large quadruple moments and vector forces¹³. Determines the transition probability of the lower electric quadrupole $B(E_2)$ of ground spin state for a first excite spin 2^+ state¹⁷:

$$(E_2 : 0^+ \rightarrow 2^+) = 5 / 16\pi e^2 Q_0^2 \dots\dots\dots 11$$

, where:

(E_2) ↑: Reduced chance of an electric quadruple transition in unit of ($e^2 b^2$).

Q_0 : the intrinsic quadruple mountain in barn (b) units.

If Q_0 Calculate elliptically charged and homogenously charged with Z charge and (b) and (a). by sign (b) relative to x-axis's, Q_0 may be¹⁸

$$Q_0 = \frac{2Z}{5}(a^2 - b^2) \dots\dots\dots 12$$

Eq 9, can be used to calculate the average, radius: $R = \frac{1}{2}(a + b)$ and $\Delta R = (b - a)$, with $\delta = \Delta R/R$, by Eq. 9, the quadruple moments is given if the deviation from sphericity is not very significant ¹⁹:

$$Q_0 = \frac{4}{5} ZR^2 \delta \dots\dots\dots 13$$

A, values of Q_0 are, determined via the equation.

Results and Discussion

Several characteristics for elemental nickel (Ni) and its isotopes, ⁵⁸Ni, ⁶⁰Ni, ⁶²Ni, ⁶⁴Ni, and ⁶⁸Ni required, have been estimated in the current study for double nuclei with masses and numbers fewer than 100 ($A > 100$). The following criteria must be met for this research:

From table 1, we observe that the lowest value of the deformation Parameter (β_2) is for the (⁶⁸Ni) equal to ($\beta_2 = 0.1530$) and the largest value of deformation parameters for (⁶²Ni) ($\beta_2 = 0.2209$). The remaining values of β_2 is, ranging, between, these (2 values). We also note that the lowest value of the deformation Parameter (δ) for the (⁶⁸Ni) is equal to ($\delta = 0.1100$) and the largest value of deformation parameters for (⁶²Ni) ($\delta = 0.1568$). We also notice from table.1, that the values deformation Parameter (δ) of the elements increase with mass numbers (58, 60, and 62) and then begin to decrease with mass numbers (64, 66, and 68). It was also noted that the values of (Q_0) range from the highest value is (1.0561) to the lowest value is (0.7778) in table 1, All of these results are represented graphically in Figs. 2,3 and 4. Also, by using the results in table 1, a relationship between the number of neutrons as a function to ($B(E_2)$) and (R_0^2), as shown in Fig. 5, 6. Also, a relationship between (δ , β_2) as a function of neutrons number was represented graphically as shown in the Fig. 7 This is because the nucleus of the (⁵⁰28Ni) is one of the nuclei with double magic numbers ($Z = 28, N = 50$), and therefore this nucleus is more stable than others as shown in the Fig. 8. Also Fig. 7, shows the correlation between the

$$\delta = 0.75Q_0 / (\langle r^2 \rangle > Z) \dots\dots\dots 14$$

The two following equations yield the semi-axes (a) and (b) ²⁰.

$$a = \sqrt{\langle r^2 \rangle (1.66 - \frac{2\delta}{0.9})} \dots\dots\dots 14$$

$$b = \sqrt{5 \langle r^2 \rangle - 2a^2} \dots\dots\dots 15$$

neutron counts and the deformation parameters. Also, through the results shown in table 1, the relationship was drawn for low electrical transition potential (B (E2) Between the present work (P W) through the results above and the theoretical study obtained from the Raman source by "single-shell Asymptotic Nilsson Model "(SSANM) and it was observed from the drawing that there is agreement between the two axes and this is due to the convergence of the values obtained in the current work with the experimental values as show in Fig. 8.

In Table 2, we note that the highest value of transition potential T(s) is equal to 8.9416×10^6 at ⁶⁶Ni where and the lowest value of T(s) is 9.6248×10^{-13} at ⁵⁸Ni and the highest value of the average half-life is 1.0390×10^{12} at ⁵⁸Ni. The lowest value is 1.1184×10^{-7} at ⁶⁶Ni, and the highest value for the gamma energy is 2033at ⁶⁸Ni. Also, the gamma energy values shown in the above table were relied on to calculate the electrical transmission probability B(E2), and through (B(E2), it was calculated the deformation coefficients (β_2, δ) through which the shape of the nuclei was determined, depending on the deformation equation on which it was based on our current study.

Observing a root mean values the squared charge values of radius $\langle r^2 \rangle^{1/2}$ in the Table 3, We discovered that values rose as mass number A increased. It was discovered that computed values $\langle r^2 \rangle^{1/2}$ present work (P w) correspond well with the experimental value of $\langle r^2 \rangle^{1/2}$, by the reference's for comparison purposes ²¹.

Table 1. the number of the mass for isotopes (A) Nickel, neutrons number (N), the energy of gamma to the first level E_γ , average nuclear radius (R_0^2), low electrical transition potential (B (E2) \uparrow in $e^2 b^2$ unit, moment Tetrapolar electrode (Q_0) in barn unit, and the parameters of the deformation (β_2, δ) to Ni.

Z	A	N	Theoretical Values			Present Work				
			E_γ (KeV)	(B(E2) \uparrow (e ² b ²) for (SSANM))	β_2 for (SSANM) (P.w)	δ	β_2	Q_0 (b)	B(E2) \uparrow (e ² b ²)	R_0^2
28	58	30	1454.28	0.016	0.0877	0.1491	0.2121	0.96320	0.0935	21.576
	60	32	1332.518	0.036	0.1286	0.1513	0.2142	1.0017	0.0998	22.069
	62	34	1172.9	0.054	0.1541	0.1568	0.2209	1.0561	0.1109	22.557
	64	36	1354.84	0.068	0.1693	0.1420	0.1991	0.9723	0.0940	23.04
	66	38	1424.8	0.079	0.1788	0.1348	0.1883	0.9384	0.0876	20.15561
	68	40	2033	0.084	0.1807	0.1100	0.1530	0.7778	0.0602	23.99

Table 2. A masses numbers (A), the neutrons numbers (N), a gamma energy of the ground level, the probability of the transition (T) and the average half-life (τ (s)) to Nickel (Ni) isotopes.

Z	A	N	E_i (keV)	E_r (keV)	$t_{1/2}$ (s)	T (s)	τ (s)
28	58	30	1454	1454.28	(667 fs) 6.67×10^{-13}	9.6248×10^{-13}	1.0390×10^{12}
	60	32	1332.501	1332.518	(0.713 ps) 7.13×10^{-13}	1.0289×10^{-12}	9.71995×10^{11}
	62	34	1172.91	1172.9	(1.45ps) 1.45×10^{-12}	2.0924×10^{-12}	4.7793×10^{11}
	64	36	1354.75	1354.84	(0.88ps) 8.8×10^{-13}	1.22698×10^{-12}	7.8750×10^{11}
	66	38	1424.8	1424.8	(54.6h) 61.965×10^5	8.9416×10^6	1.1184×10^{-7}
	68	40	2033	2033	(0.86 ms) 860	1.241×10^3	8.0581×10^{-4}

Table 3 . number's masse (A), the numbers of neutrons (N), a root mean's squares of the radius $\langle r^2 \rangle^{1/2}$, the minor & major axes (b, a) additionally, what separates them (ΔR) in two ways for isotopes of Nickel (Ni).

z	A	N	A theoretical value			present work				
			$\langle r^2 \rangle^{1/2}$ fm [21]	$\langle r^2 \rangle$ fm	$\langle r^2 \rangle^{1/2}$ fm	a (Fm)	b (fm)	ΔR_1	ΔR_2	ΔR_3
28	58	30	3.7757	17.4237	4.1741	2.3610	3.1181	0.6175	0.7572	0.9320
	60	32	3.8118	17.7345	4.2112	2.3670	3.1385	0.6338	0.7715	0.9519
	62	34	3.8399	18.0414	4.2475	2.3663	3.1685	0.6640	0.8022	0.9927
	64	36	3.8572	18.3444	4.2830	2.4057	3.1370	0.6076	0.7314	0.9043
	66	38	---	18.6438	4.3178	2.4296	3.1279	0.5830	0.6983	0.8639
	68	40	---	18.9398	4.3519	2.4879	3.0629	0.4804	0.5750	0.7090

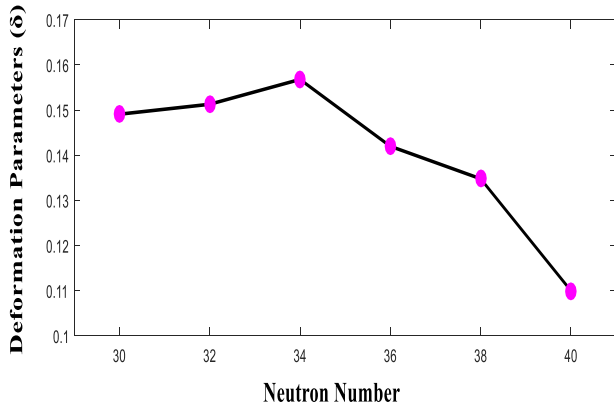


Figure 2. shows the relationship of the number of neutrons as a function of the distortion parameter δ for the element nickel (${}_{28}\text{Ni}$).

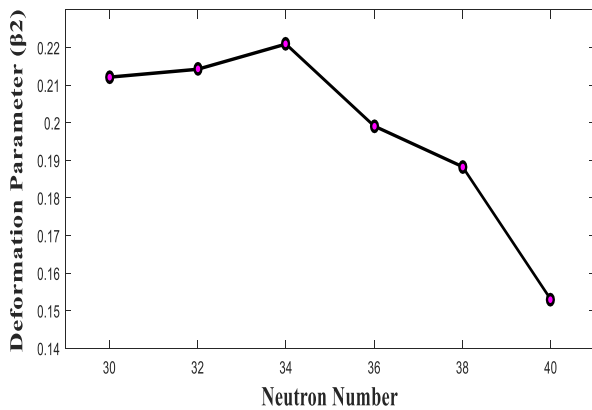


Figure 3. shows the relationship of the number of neutrons as a function of the distortion parameter β_2 for the element nickel (${}_{28}\text{Ni}$).

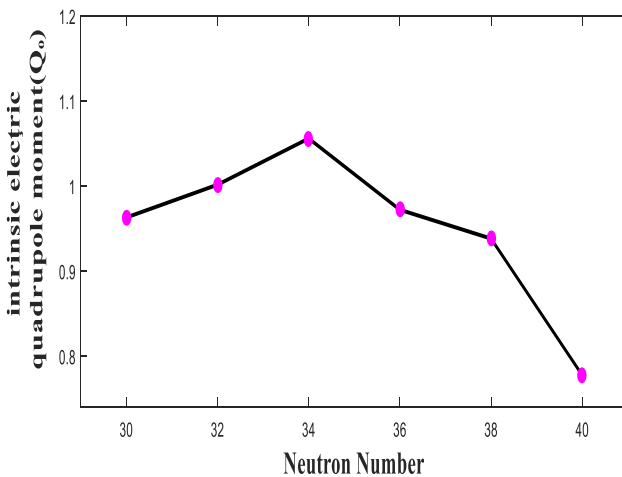


Figure 4. shows the relationship of the number of neutrons as a function of the intrinsic electric quadrupole moment (Q_0) for the element nickel (${}_{28}\text{Ni}$).

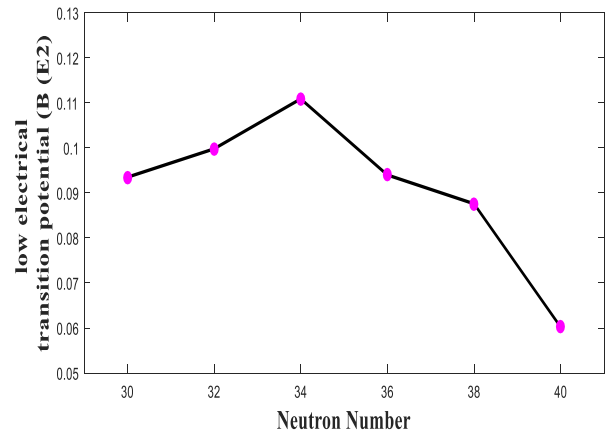


Figure 5. shows the relationship of the number of neutrons as a function of the low electrical transition potential (B (E2)) for the element nickel (${}_{28}\text{Ni}$).

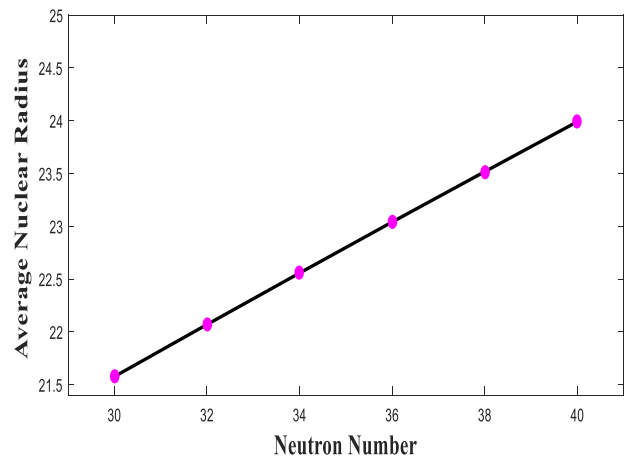


Figure 6. shows the relationship of the number of neutrons as a function of the average nuclear radius (R_0^2) for the element nickel (${}_{28}\text{Ni}$).

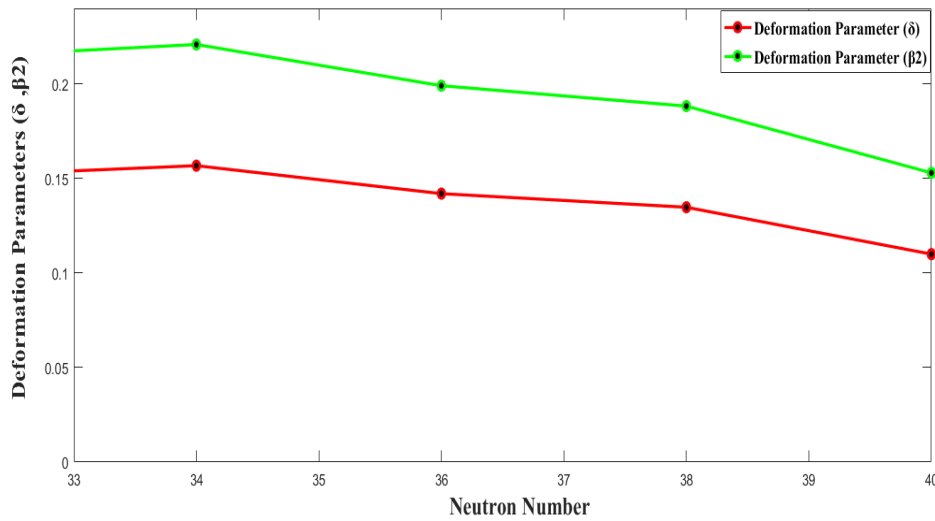


Figure 7. illustrates the relationship between the Deformation Parameters (δ, β_2) for the element nickel (^{28}Ni) as a function of the number of neutrons.

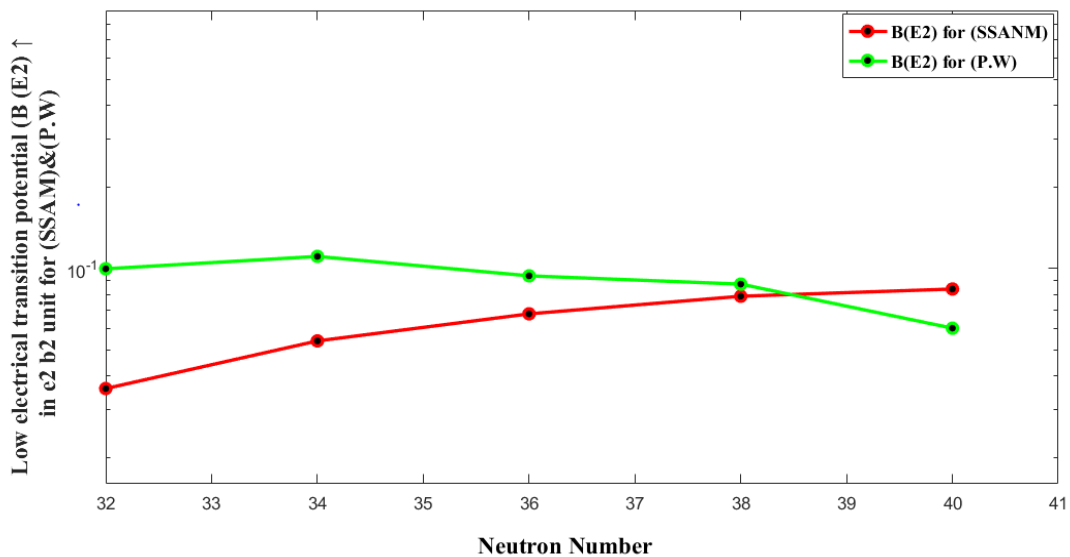


Figure 8. shows the relationship between Low electrical transition potential ($B(E2) \uparrow$) for (SSAM) & (P.W) in $e^2 b^2$ unit as a function of neutrons number for element nickel (^{28}Ni).

Conclusion

We conclude from the results obtained in the current study that the closer the nuclei, i.e. the atomic numbers or the number of neutrons for the nuclei, are to the magic numbers, the less distorted and more stable they are. Also, Nuclei show more stability and

sphericity as the mass number A increases, and the energy of the first excited state ($2+$) steadily declines as A increases, except sites close to closed shells where energy values are rather high.

Acknowledgment

Firstly, I would like to extend my sincere thanks to the participants in this research: (*Sameera Ahmed Ebrahiem*) for being the basis and first supporter in

this work, as I relied in my research on most of her scientific works of scientific research, dissertations and dissertations, in addition to (Abdallah *Ben*

Rhaim) for being an essential part in helping me and directing me to the appropriate publishing houses. In addition to relying on his scientific information in physics, I also do not forget the credit of the rest of the researchers and authors whose scientific publications I benefited from in my

research (Raman S, Nestor CW, and Tikkanen GR), as well as other things mentioned in the sources below. I also do not forget the greatest gratitude to Baghdad University, College of Science, for giving me the opportunity to publish my research in a smooth manner without complications.

Authors' Declaration

- Conflicts of Interest: None.
- We hereby confirm that all the Figures and Tables in the manuscript are ours. Furthermore, any Figures and images, that are not ours, have been included with the necessary permission for re-publication, which is attached to the manuscript.
- No animal studies are present in the manuscript.
- No human studies are present in the manuscript.
- Ethical Clearance: The project was approved by the local ethical committee at University of Sfax, Sfax, Tunisia.

Authors' Contribution Statement

S.A.E found the idea for the current study and it was the basis for choosing the topic, in addition to deducing an equation or relationship that was adopted in obtaining the results and drawings shown in the research. As for A.B.R was the first participant for S.A.E in the idea of the topic, in addition to

providing the sources that were adopted in writing. S.M.I worked on writing the research and calculating and drawing the results. All authors participated in writing the introduction and discussing the calculations and conclusions that emerged from the study.

References

1. Hameed BS. The Nuclear Structure for Exotic Neutron-Rich of ^{42,43,45,47}K Nuclei. Baghdad Sci J. 2020 Apr; 13(1): 146-9. <http://dx.doi.org/10.21123/bsj.2016.13.1.0146>
2. Akkoyun S, Bayram T, Kara SO. A study on estimation of electric quadrupole transition probability in nuclei. J Nucl Sci. 2015 Jan 5; 2(1): 7-10. <https://doi.org/10.1501/nuclear.0000000009>.
3. Kumar R, Bhuyan M, Jain D, Carlson BV. Theoretical Description of Low-Energy Nuclear Fusion. In Nuclear Structure Physics. CRC Press. 2020 Oct ; 21: 121-144. <https://doi.org/10.1201/9780429288647-5>
4. Raheem EM, Hasan AA, Alwan IH. Study of ground state properties of some Ni-Isotopes using Skyrme-Hartree-Fock method. Iraqi J Phys. 2019 Jul 1; 17(42): 1-2. <https://doi.org/10.20723/ijp.17.42.1-12>
5. Ali AH, Hassoon SO, Tafash HT. Calculations of Quadrupole Deformation Parameters for Nuclei in fp shell. In J Phys Conf Ser. 2019 Feb; 1178: 012010. <http://dx.doi.org/10.1088/1742-6596/1178/1/012010>
6. Heyde KL, Heyde KL. The nuclear shell model. Springer Berlin Heidelberg; 1994. <https://doi.org/10.1007/978-3-642-97203-4>
7. Mahmood PF. Ground State Properties of Even-Even 30– 92Ca Isotopes Using HFB Theory. Kirkuk J. Sci. 2024 Mar 1;19(1):43-50. <https://doi.org/10.32894/KUJSS.2024.146573.1136>
8. Raman S, Nestor Jr CW, Tikkanen P. Transition probability from the ground to the first-excited 2+ state of even-even nuclides. At Data Nucl. Data Tables. 2001 May 1; 78(1): 1-28. <http://dx.doi.org/10.1006/adnd.2001.0858>
9. Jakubus M, Graczyk M. Availability of nickel in soil evaluated by various chemical extractants and plant accumulation. Agronomy. 2020 Nov 17; 10(11): 1-16. <http://dx.doi.org/10.3390/agronomy10111805>
10. Adamu A. A New Measurement of Nuclear Radius from the study of β^+ -Decay Energy of Finite-Sized Nuclei. J Rad Nucl Appl. 2021; 6(1): 45-49. <http://dx.doi.org/10.18576/jrna/060108>
11. Ibrahim SM, Ebrahiem SA, Rhaim AB. Calculating of the Electric Transition Forces and Radii of Even-even-Nuclei of Cadmium (100-12448Cd) Cd Isotopes. Ibn al-Haitham J Pure Appl Sci. 2024 Apr 20; 37(2): 217-25. <https://doi.org/10.30526/37.2.3468>
12. Majeed FA, Obaid SM. Nuclear structure study of ²², ²⁴Ne and ²⁴Mg nuclei. Revista mexicana de física. 2019 Apr; 65(2): 159-67. <http://dx.doi.org/10.31349/RevMexFis.65.159>

13. Alzubadi AA, Dakhil ZA, Aluboodi ST. Microscopic Study of Nuclear Structure for Some Zr-isotopes Using Skyrme-Hartree-Fock-Method. Journal of Nuclear and Particle Physics. 2014;4(6):155-163.<https://doi.10.5923/j.jnpp.20140406.01>
14. Kenichi Yoshida. Suddenly shortened half-lives beyond ^{78}Ni : $N=50$ magic number and high-energy nonunique first-forbidden transitions. Phys Rev C. 2019; 100(2): 1-9.<https://doi.10.1103/PhysRevC.100.024316>
15. Angeli I, Marinova KP. Table of experimental nuclear ground state charge radii: An update. At Data Nucl. Data Tables. 2013 Jan 1; 99(1): 69-95.<http://dx.doi.org/10.1016/j.adt.2011.12.006>
16. Hameed BS, Rejah BK. Study the Nuclear Structure of Some Cobalt Isotopes. Baghdad Sci J. 2022 Dec 5;19(6 (Suppl.)):1566-1571.<https://dx.doi.org/10.21123/bsj.2022.7537>
17. Yakut H, Tabar E, Kemah E, Hoşgör G. Microscopic calculation of the electromagnetic dipole strength for $^{239,243}\text{Pu}$ isotopes. J. Phys. G: Nucl. Part. Phys. 2022 Dec 13;50(1):015104.<https://doi.10.1088/1361-6471/aca3bf>
18. Taqi A, Mahmood P. Ground State Properties and Shape Transition of Even-Even ^{76}Os , ^{78}Pt , ^{80}Hg and ^{82}Pb Isotopes in the Framework of Skyrme Hartree-Fock-Bogoliubov Theory. Arab J. Nucl. Sci. Appl. 2022 Jan 1;55(1):34-44. <https://doi.10.21608/ajnsa.2021.70418.1461>
19. Ye W, Qian Y, Ren Z. Refining the nuclear mass model via the α decay energy. Phys. Rev. C. 2021 Dec;104(6):064308.<https://doi.org/10.1103/PhysRevC.104.064308>
20. Taha H, Ibrahim K, Rahman MM, Henry DJ, Yin CY, Veder JP, et al. Sol-gel derived ITO-based bi-layer and tri-layer thin film coatings for organic solar cells applications. Appl Surf Sci. 2020 Nov 15; 530: 147164. <https://doi.org/10.1016/j.apsusc.2020.147164>
21. Abd Ali FS, Mahdi KH, Jawad EA. Humidity effect on diffusion and length coefficient of radon in soil and building materials. Energy Procedia. 2019 Jan 1; 157: 384-392.<https://doi.org/10.1016/j.egypro.2018.11.203>

حساب انصاف اقطار نظائر النيكل الزوجية-زوجية

صفا محسن ابراهيم¹، عبد الله بن رحيم¹، سميرة احمد ابراهيم²

¹قسم الفيزياء، كلية العلوم، جامعة صفاقس، صفاقس، تونس.
²قسم الفيزياء، كلية التربية ابن الهيثم للعلوم الصرفة، جامعة بغداد، بغداد، العراق.

الخلاصة

ركز البحث الحالي على دراسة اشكال النوى زوجية-زوجية لنظائر النيكل والتي لها اعداد كتلية اقل من 100 حيث تم دراسة معاملات التشوه (β_2) والمشتقة من احتمالية الانتقال الكهربائي $B(E2)$ وكذلك معاملات التشوه (δ) المشتقة من العزم الكهربائي رباعي القطب Q_0 وتم دراسة متوسط مربع نصف القطر $\langle r^2 \rangle^{1/2}$ بالإضافة الى المحاور الصغرى والكبرى (a,b) والفرق بينهما باستخدام معادلات نموذج القشرة المشوهة حيث تم ملاحظة الاختلاف في اشكال النوى لنظائر النيكل من خلال رسم اشكال ثلاثية وثنائية الابعاد ومن خلال النتائج التي تم الحصول عليها لوحظ ان النظائر الاكثر تشوها وميلا لعدم الاستقرار واتخاذ الشكل الاهليلجي هي النظائر ذات الاعداد النيوترونات والبروتونات البعيدة عن الاعداد السحرية وكذلك تم مقارنة النتائج التي تم الحصول عليها مع النتائج النظرية من المصدر (SSANM) حيث لوحظ انها مقاربة لها نوعا ما.

الكلمات المفتاحية: احتمالية الانتقال ($B(E2: 0+ \rightarrow 2+)$ ، عزم الكهربائي رباعي القطب Q_0 ، معاملات التشوه (δ , β_2)، متوسط مربع توزيع الشحنة، نصف القطر $\langle r^2 \rangle$.