

Using the Elzaki decomposition method to solve nonlinear fractional differential equations with the Caputo-Fabrizio fractional operator

Mohammed Abdulshareef Hussein  

Scientific Research Center, AL-Ayen University, Thi-Qar, Iraq

Received 09/04/2022, Revised 24/02/2023, Accepted 26/02/2023, Published Online First 20/08/2023,
Published 01/03/2024



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Abstract

The techniques of fractional calculus are applied successfully in many branches of science and engineering, one of the techniques is the Elzaki Adomian decomposition method (EADM), which researchers did not study with the fractional derivative of Caputo Fabrizio. This work aims to study the Elzaki Adomian decomposition method (EADM) to solve fractional differential equations with the Caputo-Fabrizio derivative. We presented the algorithm of this method with the CF operator and discussed its convergence by using the method of the Cauchy series then, the method has applied to solve Burger, heat-like, and, coupled Burger equations with the Caputo -Fabrizio operator. To conclude the method was convergent and effective for solving this type of fractional differential equations.

Keywords: Burger equation, Caputo-Fabrizio fractional operator, Elzaki decomposition method, Heat-like equation.

Introduction

Fractional calculus (that is, calculus of integrals and derivatives of any arbitrary real or complex order) has grown in popularity and relevance over the last three decades, owing to its proven applications in a wide range of apparently disparate domains of science and engineering. It does, in fact, give some potentially valuable methods for solving differential and integral equations, as well as a variety of other problems requiring mathematical physics special functions, as well as their extensions and generalizations in one or more variables¹⁻⁴.

In the past decade, Caputo and Fabrizio introduced a new fractional differential operator many researchers studied this operator and researchers are still interested in this operator because of its importance, as some studies have applied methods of approximate solutions to equations that include this fractional operator^{5,6}.

The requirement for a model that represents the behavior of classic viscoelastic materials, thermal media, electromagnetic systems, and so on has piqued the attention of researchers in this unique technique. Plasticity, fatigue, and damage, as well as electromagnetic hysteresis, appear to benefit most from the original concept of fractional derivative. When these effects are missing, it appears that the new fractional derivative is more appropriate⁷.

Despite the importance of integrative transformations in solving differential equations, most studies have focused on some of those transformations and neglected other transformations. Among these transformations is the Elzaki transformation presented by researcher Tariq Elzaki in recent years. Therefore, in this study, The Elzaki decomposition method will be presented for solving differential equations that include the fractional operator Caputo-Fabrizio⁸⁻¹¹.

Fractional calculus preliminaries

Definition 1: ¹² The fractional derivative with the Caputo-Fabrizio operator for $0 < a \leq 1$ is defined as:

$${}^{CF}D_t^a u(t) = \frac{\varepsilon(a)}{1-a} \int_0^t \exp\left[-\frac{a(t-s)}{1-a}\right] u'(s) ds, \quad 1$$

where $u \in \mathcal{H}^1(z_1, z_2)$, $z_1 < z_2$, $u'(s)$ is the derivative of u , and $\varepsilon(a)$ is a normalization function such that $\varepsilon(0) = \varepsilon(1) = 1$.

The operator's fundamental attributes are as follows ^{13,14}

1. ${}^{CF}D_t^a u(t) = u(t)$, where $a = 0$.
2. ${}^{CF}D_t^a [u(t) + v(t)] = {}^{CF}D_t^a u(t) + {}^{CF}D_t^a v(t)$, where $u, v \in \mathcal{H}^1(z_1, z_2)$
3. ${}^{CF}D_t^a (c) = 0$, c is constant.

Definition 2: ¹⁵ Over a set of functions \mathcal{A} , the Elzaki transform is defined,

$$\mathcal{A} = \left\{ u(t) \left| \begin{array}{l} \exists \mathcal{M}, \tau_1, \tau_2 > 0, |u(t)| < \mathcal{M} e^{\frac{|t|}{\tau_j}} \\ \text{if } t \in (-1)^j \times [0, \infty) \end{array} \right. \right\},$$

by the following formula

$$E\{u(t)\} = w \int_0^\infty u(t) e^{-\frac{t}{w}} dt, \quad w \in (\tau_1, \tau_2), \quad 2$$

where w is a parameter of Elzaki transform.

Elzaki transform for some functions ¹⁵⁻¹⁹:

1. $E\{1\} = w^2$.
2. $E\{t\} = w^3$.
3. $E\{t^n\} = n! w^{n+2}$.
4. $E\{e^{at}\} = \frac{w^2}{1-aw}$.
5. $E\{\sin(at)\} = \frac{aw^3}{1+a^2w^2}$.
6. $E\{\cos(at)\} = \frac{w^2}{1+a^2w^2}$.

Lemma 1: The Elzaki transform for the Caputo-Fabrizio fractional operator is defined as follows if

$$0 < a \leq 1.$$

$$E\{{}^{CF}D_t^a u(t)\} = \frac{[E\{u(t)\} - w^2 u(0)]}{1-a+aw}, \quad 3$$

Proof:

$$\begin{aligned} & E\{{}^{CF}D_t^a u(t)\} \\ &= E\left\{ \frac{1}{1-a} \int_0^t u'(s) \exp\left[-\frac{a(t-s)}{1-a}\right] ds \right\}, \\ &= \frac{1}{1-a} E\left\{ \int_0^t u'(s) \exp\left[-\frac{a(t-s)}{1-a}\right] ds \right\}, \\ &= \frac{w^{-1}}{1-a} E\{u'(t)\} E\left\{ \exp\left[-\frac{at}{1-a}\right] \right\}, \\ &= \frac{w^{-1}}{1-a} \left[\frac{1}{w} E\{u(t)\} - w u(0) \right] \left[\frac{w^2(1-a)}{1-a+aw} \right], \\ &= \frac{[E\{u(t)\} - w^2 u(0)]}{1-a+aw}. \end{aligned}$$

Analysis of EDM

In the Caputo-Fabrizio operator sense, consider the following nonlinear partial differential equation:

$${}^{CF}D_t^a u(x, t) + \mathcal{R}[u] + \mathcal{N}[u] = \mathcal{G}(x, t), \quad 4$$

$$\text{with initial condition } u(x, 0) = u_0(x), \quad 5$$

where ${}^{CF}D_t^a u(x, t)$ is Caputo-Fabrizio operator of $u(x, t)$, a linear operator is \mathcal{R} , a nonlinear operator is \mathcal{N} and a source term is \mathcal{G} .

Applying the Elzaki transform to both sides of the Eq.4,

$$E\{{}^{CF}D_t^a u(x, t) + \mathcal{R}[u] + \mathcal{N}[u]\} = E\{\mathcal{G}\}, \quad 6$$

from Lemma 1 and Eq.5,

$$\frac{[E\{u(t)\} - w^2 u(0)]}{1-a+aw} = E\{\mathcal{G} - \mathcal{R}[u] - \mathcal{N}[u]\}, \quad 7$$

or

$$\begin{aligned} E\{u\} &= w^2 u_0(x) + (1-a+aw) E\{\mathcal{G}\} \\ &\quad - (1-a+aw) E\{\mathcal{R}[u]\} \\ &\quad - (1-a+aw) E\{\mathcal{N}[u]\}, \quad 8 \end{aligned}$$

by using the Elzaki transform's inverse to Eq.8,

$$\begin{aligned} u(x, t) &= E^{-1}\{w^2 u_0(x)\} \\ &\quad + E^{-1}\{(1-a+aw) E\{\mathcal{G}\}\} \\ &\quad - E^{-1}\{(1-a+aw) E\{\mathcal{R}[u]\}\} \\ &\quad - E^{-1}\{(1-a+aw) E\{\mathcal{N}[u]\}\}. \quad 9 \end{aligned}$$

Suppose that $u(x, t)$ is a solution of Eq.9, which it expressed as

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t), \quad 10$$

the nonlinear term can be decomposed as

$$\mathcal{N}[u(x, t)] = \sum_{n=0}^{\infty} \mathcal{A}_n, \quad 11$$

where,

$$\mathcal{A}_n = \frac{1}{n!} \frac{\partial^n}{\partial \zeta^n} [\mathcal{N}(\sum_{i=0}^n \zeta^i u_i(x, t))]_{\zeta=0}, \quad 12$$

where, $n = 0, 1, 2, \dots$

Substituting Eq.10 and Eq.11 into Eq.9 gives us the result that

$$\begin{aligned} \sum_{n=0}^{\infty} u_n = & E^{-1}\{w^2 u_0(x)\} + E^{-1}\{(1-a+aw)E\{g\}\} \\ & - E^{-1}\{(1-a+aw)E\{\mathcal{R}[\sum_{n=0}^{\infty} u_n]\}\} \\ & - E^{-1}\{(1-a+aw)E\{\sum_{n=0}^{\infty} \mathcal{A}_n\}\}. \end{aligned} \quad 13$$

When the left and right sides of Eq.13 are compared,

$$\begin{aligned} u_0 &= u_0(x) + E^{-1}\{(1-a+aw)E\{g\}\}, \\ u_1 &= -E^{-1}\{(1-a+aw)E\{\mathcal{R}[u_0]\}\} \\ & \quad - E^{-1}\{(1-a+aw)E\{\mathcal{A}_0\}\}, \\ u_2 &= -E^{-1}\{(1-a+aw)E\{\mathcal{R}[u_1(x, t)]\}\} \\ & \quad - E^{-1}\{(1-a+aw)E\{\mathcal{A}_1\}\}. \end{aligned} \quad 14$$

In its general form, a recursive relation is

$$\begin{aligned} u_0 &= u_0(x) + E^{-1}\{(1-a+aw)E\{g\}\}, \\ u_{n+1} &= -E^{-1}\{(1-a+aw)E\{\mathcal{R}[u_n]\}\} \\ & \quad - E^{-1}\{(1-a+aw)E\{\mathcal{A}_n\}\}. \end{aligned} \quad 15$$

The approximate solution is given by

$$u(x, t) = u_0 + u_1 + u_2 + \dots = \sum_{i=0}^{\infty} u_i. \quad 16$$

Convergence analysis

This section discusses the convergence of Elzaki decomposition method to the exact solution of fractional differential equations as well as the estimated error resulting from the approximate solutions.

Theorem 1: Suppose that B is a Banach space, $\sum_{i=0}^{\infty} u_i$ in Eq.16 is convergence to $\delta \in B$ if $\exists(0 \leq \alpha < 1)$, s.t. $\forall \eta \in \mathbb{N} \Rightarrow \|u_\eta\| \leq \alpha \|u_{\eta-1}\|$.

Proof: The sequence of partial sums is defined as $\{\delta_\eta\}_{\eta=0}^{\infty}$

$$\begin{aligned} \delta_0 &= u_0 \\ \delta_1 &= u_0 + u_1 \\ \delta_2 &= u_0 + u_1 + u_2 \\ &\vdots \\ \delta_\eta &= u_0 + u_1 + \dots + u_\eta, \end{aligned}$$

now, it is necessary to prove that, $\{\delta_\eta\}_{\eta=0}^{\infty}$ is a Cauchy series in Banach space,

$$\begin{aligned} \|\delta_{\eta+1} - \delta_\eta\| &= \|\sum_{i=0}^{\eta+1} u_i - \sum_{i=0}^{\eta} u_i\| = \|u_{\eta+1}\| \leq \\ &\alpha \|u_\eta\| \leq \dots \leq \alpha^{\eta+1} \|u_0\|, \end{aligned} \quad 17$$

for all $\eta, m \in \mathbb{N}$ as $\eta \geq m$,

$$\begin{aligned} \|\delta_\eta - \delta_m\| &= \|(\delta_\eta - \delta_{\eta-1}) + (\delta_{\eta-1} - \delta_{\eta-2}) + \dots \\ &\quad + (\delta_{m+1} - \delta_m)\| \\ &\leq \|\delta_\eta - \delta_{\eta-1}\| + \|\delta_{\eta-1} - \delta_{\eta-2}\| + \dots \\ &\quad + \|\delta_{m+1} - \delta_m\| \\ &\leq \alpha^\eta \|u_0\| + \alpha^{\eta-1} \|u_0\| + \dots + \\ &\quad \alpha^{m+1} \|u_0\| \\ &= \alpha^{m+1} \|u_0\| (\alpha^{\eta-m-1} + \alpha^{\eta-m-2} + \dots + \alpha) \\ &= \frac{1-\alpha^{\eta-m}}{1-\alpha} \alpha^{m+1} \|u_0\|. \end{aligned} \quad 18$$

Since $(\alpha^{\eta-m-1} + \alpha^{\eta-m-2} + \dots + \alpha)$ is a geometric series and $0 \leq \alpha < 1$, then,

$$\lim_{\eta, m \rightarrow \infty} \|\delta_\eta - \delta_m\| = 0,$$

thus, $\{\delta_\eta\}_{\eta=0}^{\infty}$ is Cauchy sequence in Banach space B , therefore produces that the series solution $u = \sum_{i=0}^{\infty} u_i$, defined in Eq.16 converges.

Theorem 2: Suppose that the series solution $\sum_{i=0}^{\infty} u_i$ in Eq.16 is convergent to the solution $u(x, t)$. If $\sum_{i=0}^{\infty} u_i$ is used as an approximation to the solution $u(x, t)$ of Eq.5 then the maximum error, $E_m(x, t)$ is estimated as

$$E_m(x, t) \leq \frac{1}{1-\alpha} \alpha^{m+1} \|u_0\|.$$

Proof: From Theorem 1, inequality 18

$$\|\delta_\eta - \delta_m\| \leq \frac{1-\alpha^{\eta-m}}{1-\alpha} \alpha^{m+1} \|u_0\|,$$

for $\eta \geq m$, now, as $\eta \rightarrow \infty$ then $\delta_\eta \rightarrow u(x, t)$ so,

$$\|u(x, t) - \sum_{i=0}^{\infty} u_i\| \leq \frac{1-\alpha^{\eta-m}}{1-\alpha} \alpha^{m+1} \|u_0\|. \quad 19$$

Also, since $0 \leq \alpha < 1$ produces $(1 - \alpha^{\eta-m}) < 1$.

Therefore the above inequality 19 becomes

$$E_m(x, t) \leq \frac{1}{1-\alpha} \alpha^{m+1} \|u_0\|. \quad 20$$

Illustrative examples

Example 1: Consider the following, Burger's equation in the Caputo-Fabrizio sense is nonlinear.

$${}^{CF}_0 D_t^a u + u \frac{\partial u}{\partial x} = \frac{\partial^2 u}{\partial x^2}, \quad 0 < a \leq 1, \quad 21$$

subject to the initial condition $u(x, 0) = x$. 22

By taking the Elzaki transform to both sides of Eq.21,

$$E\{{}^{CF}_0 D_t^a u\} = w^2 x + E\left\{\frac{\partial^2 u}{\partial x^2}\right\} - E\left\{u \frac{\partial u}{\partial x}\right\}, \quad 23$$

using the inverse Elzaki transform to both sides of Eq.23,

$$u = x + E^{-1}\left\{(1 - a + aw) \left[E\left\{\frac{\partial^2 \sum_{n=0}^{\infty} u_n}{\partial x^2}\right\}\right] - E\left\{\sum_{n=0}^{\infty} \mathcal{A}_n\right\}\right\}. \quad 24$$

As a result, the approximate solution may be derived using Eq.15

$$u_0(x, t) = x,$$

$$u_1(x, t) = E^{-1}\left\{(1 - a + aw) \left[E\{\mathcal{A}_0\} - E\left\{\frac{\partial^2 u_0}{\partial x^2}\right\}\right]\right\},$$

$$u_2(x, t) = E^{-1}\left\{(1 - a + aw) \left[E\{\mathcal{A}_1\} - E\left\{\frac{\partial^2 u_1}{\partial x^2}\right\}\right]\right\},$$

by the above algorithms and after simple steps,

$$u_0 = x,$$

$$u_1 = -x(1 - a + at),$$

$$u_2 = x \begin{pmatrix} 2a^2 - 4a + a^2 t^2 \\ -4a^2 t + 4at + 2 \end{pmatrix}, \quad 25$$

⋮

and so on.

Therefore, the series solution $u(x, t)$ of Eq.21 is given by

$$u(x, t) = x - x(1 - a + at) + x \begin{pmatrix} 2a^2 - 4a + a^2 t^2 \\ -4a^2 t + 4at + 2 \end{pmatrix} - \dots \quad 26$$

If it was $a \rightarrow 1$ in Eq.26 then, the exact solution is

$$u(x, t) = x - xt + xt^2 - \dots, \\ = x \sum_{k=0}^{\infty} (-t)^k = \frac{x}{1+t}. \quad 27$$

Table 1 displays the values of the exact and approximate solutions of Eq.21 at the different a values, also includes the absolute error of the approximate solutions for the exact solutions. Fig. 1 shows graphs of the exact and approximate solutions of Eq.21 at the different a values between 0 and 1.

Table 1. The approximate and exact solution of Burger's equation with fractional operator CF, u_1 is app. sol. of u at $a = 0.9$, u_2 is app. sol. of u at $a = 1$, u_1 is exact sol. of u .

x	t	u_1	u_2	u_3	$ u_1 - u_3 $	$ u_2 - u_3 $
0.2500	0.2500	0.2089	0.2031	0.2000	0.0089	0.0031
0.7500	0.2500	0.6267	0.6094	0.6000	0.0267	0.0094
0.2500	0.5000	0.2131	0.1875	0.1667	0.0465	0.0208
0.5000	0.5000	0.4262	0.3750	0.3333	0.0929	0.0417
0.7500	0.5000	0.6394	0.5625	0.5000	0.1394	0.0625
0.2500	0.7500	0.2427	0.2031	0.1429	0.0998	0.0603
0.5000	0.7500	0.4853	0.4062	0.2857	0.1996	0.1205
0.7500	0.7500	0.7280	0.6094	0.4286	0.2994	0.1808

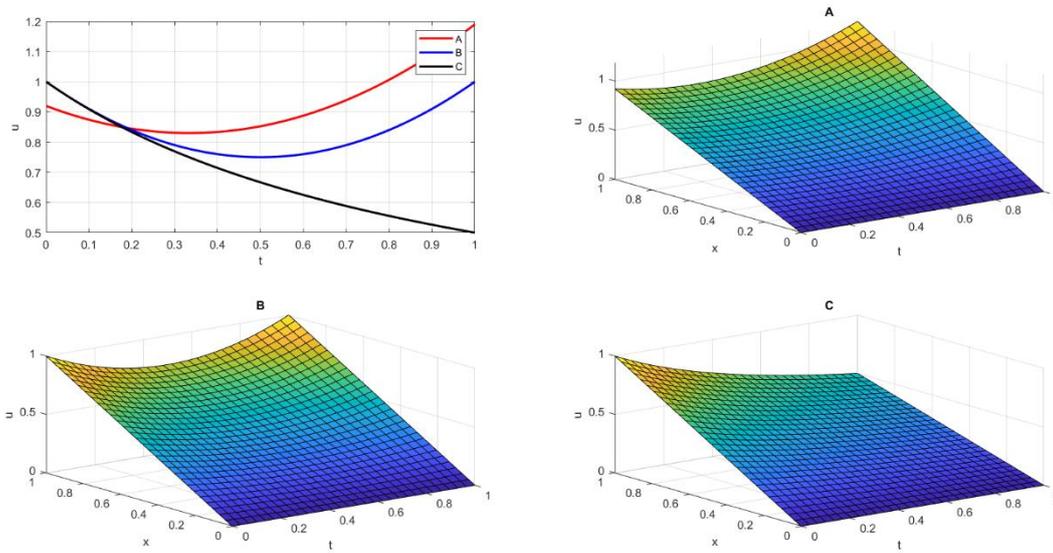


Figure 1. The graphs of the approximate and the exact solutions among different values of x and t in case (3D) and fixed x in case (2D) when $a = 0.9, 1$ for nonlinear Burger's equation in the CF fractional operator, (A) is the app. sol. of u at $a = 0.9$, (B) is the app. sol. of u at $a = 1$, (C) is the exact sol. of u .

Example 2: Consider the following heat-like equation in the Caputo-Fabrizio sense

$${}^{CF}_0 D_t^a u(x, y, t) = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}, \quad (28)$$

where $0 \leq x, y \leq 2\pi, t > 0$, with the initial condition

$$u(x, y, 0) = \sin(x) \sin(y), \quad (29)$$

applying the Elzaki transform differentiation property on Eq.28

$$E\{{}^{CF}_0 D_t^a u\} = w^2 \sin(x) \sin(y) + (1 - a + aw) \left[\begin{array}{l} E\left\{\frac{\partial^2 u}{\partial x^2}\right\} \\ + E\left\{\frac{\partial^2 u}{\partial y^2}\right\} \end{array} \right], \quad (30)$$

by using the inverse Elzaki transform to both sides of Eq.30,
 $u = \sin(x) \sin(y)$

$$+ E^{-1} \left\{ (1 - a + aw) \left[\begin{array}{l} E\left\{\frac{\partial^2 u}{\partial x^2}\right\} \\ + E\left\{\frac{\partial^2 u}{\partial y^2}\right\} \end{array} \right] \right\}, \quad (31)$$

as a result, the approximate solution may be derived using Eq.15,

$$u_0 = \sin(x) \sin(y),$$

$$u_n = E^{-1} \left\{ (1 - a + aw) \left[\begin{array}{l} E\left\{\frac{\partial^2 u_{n-1}}{\partial x^2}\right\} \\ + E\left\{\frac{\partial^2 u_{n-1}}{\partial y^2}\right\} \end{array} \right] \right\}. \quad (32)$$

Hence, from Eq 29 and Eq 32, the components give as follows:

$$u_0 = \sin(x) \sin(y),$$

$$u_1 = E^{-1} \left\{ (1 - a + aw) \left[\begin{array}{l} E\left\{\frac{\partial^2 u_0}{\partial x^2}\right\} \\ + E\left\{\frac{\partial^2 u_0}{\partial y^2}\right\} \end{array} \right] \right\} = -2 \sin(x) \sin(y) (1 - a + at)$$

$$u_2 = E^{-1} \left\{ (1 - a + aw) \left[\begin{array}{l} E\left\{\frac{\partial^2 u_0}{\partial x^2}\right\} \\ + E\left\{\frac{\partial^2 u_0}{\partial y^2}\right\} \end{array} \right] \right\} = 4 \sin(x) \sin(y) \left[(1 - 2a + a^2) + (2a - 2a^2)t + \frac{1}{2} a^2 t^2 \right]. \quad (33)$$

and so on.

Therefore, the approximate solution of $u(x, y, t)$ of Eq.28 is given by

$$u = \sin(x) \sin(y) - 2 \sin(x) \sin(y) (1 - a + at) + 4 \sin(x) \sin(y) \left[\begin{array}{l} (1 - 2a + a^2) \\ + (2a - 2a^2)t + \frac{1}{2} a^2 t^2 \end{array} \right] + \dots \quad (34)$$

If it was $a \rightarrow 1$ in Eq.34 then, the exact solution is,

$$u = \sin(x) \sin(y) \left(1 - 2t + \frac{(2t)^2}{2!} - \dots\right),$$

$$= \sin(x) \sin(y) e^{-2t}. \quad 35$$

Table 2 displays the values of the exact and approximate solutions of Eq.28 at the different a values, also includes the absolute error of the

approximate solutions for the exact solutions. Fig. 2 shows graphs of the exact and approximate solutions of Eq.28 at the different a values between 0 and 1.

Table 2. The approximate and exact solution of heat-like equation with fractional operator CF, u_1 is app. sol. of u at $a = 0.9$, u_2 is app. sol. of u at $a = 1$, u_3 is exact sol. of u .

x	y	t	u_1	u_2	u_3	$ u_1 - u_3 $	$ u_2 - u_3 $
0.2500	0.2500	0.2500	0.0411	0.0383	0.0371	0.0040	0.0011
0.5000	0.2500	0.2500	0.0796	0.0741	0.0719	0.0077	0.0022
0.7500	0.2500	0.2500	0.1132	0.1054	0.1023	0.0109	0.0031
0.2500	0.5000	0.2500	0.0796	0.0741	0.0719	0.0077	0.0022
0.5000	0.5000	0.2500	0.1543	0.1437	0.1394	0.0149	0.0042
0.7500	0.5000	0.2500	0.2194	0.2042	0.1982	0.0211	0.0060
0.2500	0.7500	0.2500	0.1132	0.1054	0.1023	0.0109	0.0031
0.5000	0.7500	0.2500	0.2194	0.2042	0.1982	0.0211	0.0060
0.7500	0.7500	0.2500	0.3119	0.2904	0.2818	0.0301	0.0086
0.2500	0.2500	0.5000	0.0432	0.0306	0.0225	0.0206	0.0081
0.5000	0.2500	0.5000	0.0836	0.0593	0.0436	0.0400	0.0157
0.7500	0.2500	0.5000	0.1189	0.0843	0.0620	0.0569	0.0223
0.2500	0.5000	0.5000	0.0836	0.0593	0.0436	0.0400	0.0157
0.5000	0.5000	0.5000	0.1620	0.1149	0.0846	0.0775	0.0304
0.7500	0.5000	0.5000	0.2304	0.1634	0.1202	0.1102	0.0432
0.2500	0.7500	0.5000	0.1189	0.0843	0.0620	0.0569	0.0223
0.5000	0.7500	0.5000	0.2304	0.1634	0.1202	0.1102	0.0432
0.7500	0.7500	0.5000	0.3276	0.2323	0.1709	0.1566	0.0614
0.2500	0.2500	0.7500	0.0576	0.0383	0.0137	0.0440	0.0246
0.5000	0.2500	0.7500	0.1116	0.0741	0.0265	0.0852	0.0477
0.7500	0.2500	0.7500	0.1587	0.1054	0.0376	0.1211	0.0678
0.2500	0.5000	0.7500	0.1116	0.0741	0.0265	0.0852	0.0477
0.5000	0.5000	0.7500	0.2163	0.1437	0.0513	0.1651	0.0924
0.7500	0.5000	0.7500	0.3076	0.2042	0.0729	0.2347	0.1313
0.2500	0.7500	0.7500	0.1587	0.1054	0.0376	0.1211	0.0678
0.5000	0.7500	0.7500	0.3076	0.2042	0.0729	0.2347	0.1313
0.7500	0.7500	0.7500	0.4373	0.2904	0.1037	0.3337	0.1867

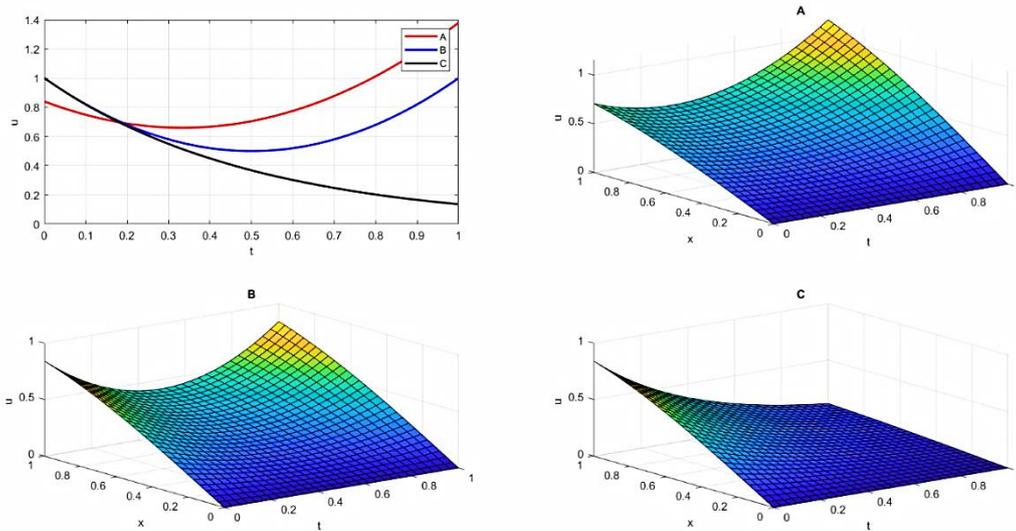


Figure 2. The graphs of the approximate and the exact solutions of heat-like equation among different values of x and t when y is fixed in case (3D) and x, y are fixed in case (2D) when $a = 0.9, 1$ for nonlinear Burger's equation in the CF fractional operator, (A) is the app. sol. of u at $a = 0.9$, (B) is the app. sol. of u at $a = 1$, (C) is the exact sol. of u .

Example 3: Consider the Caputo-Fabrizio operator's nonlinear system of time-fractional differential equations:

$${}^{\text{CF}}D_t^a u(x, t) - u_{xx} - 2uu_x + (uv)_x = 0,$$

$${}^{\text{CF}}D_t^a v(x, t) - v_{xx} - 2vv_x + (uv)_x = 0, \quad 36$$

where $0 < a, b \leq 1$ and the initial conditions are

$$\begin{aligned} u(x, 0) &= \sin(x), \\ v(x, 0) &= \sin(x). \end{aligned} \quad 37$$

Taking the Elzaki transform on both sides of Eq.36,

$$\begin{aligned} E\{ {}^{\text{CF}}D_t^a u(x, t) \} &= w^2 \sin(x) \\ &+ (1 - a + aw) E \left\{ \frac{\partial^2 u}{\partial x^2} + 2u \frac{\partial u}{\partial x} - \frac{\partial}{\partial x} (uv) \right\}, \end{aligned}$$

$$\begin{aligned} E\{ {}^{\text{CF}}D_t^a v(x, t) \} &= w^2 \sin(x) \\ &+ (1 - a + aw) E \left\{ \frac{\partial^2 v}{\partial x^2} + 2v \frac{\partial v}{\partial x} - \frac{\partial}{\partial x} (uv) \right\}, \end{aligned} \quad 38$$

when taking the inverse of the Elzaki transform to both sides of Eq.38,

$$\begin{aligned} u(x, t) &= \sin(x) \\ &+ E^{-1} \left\{ (1 - a + aw) E \left\{ \begin{aligned} &\frac{\partial^2 u}{\partial x^2} \\ &+ 2u \frac{\partial u}{\partial x} \\ &- \frac{\partial}{\partial x} (uv) \end{aligned} \right\} \right\}, \end{aligned}$$

$$v(x, t) = \sin(x)$$

$$+ E^{-1} \left\{ (1 - b + bw) E \left\{ \begin{aligned} &\frac{\partial^2 v}{\partial x^2} \\ &+ 2v \frac{\partial v}{\partial x} \\ &- \frac{\partial}{\partial x} (uv) \end{aligned} \right\} \right\}, \quad 39$$

the solution is now represented as an infinite series as seen below

$$\begin{aligned} u(x, t) &= \sum_{n=0}^{\infty} u_n(x, t), \\ v(x, t) &= \sum_{n=0}^{\infty} v_n(x, t), \end{aligned} \quad 40$$

and the nonlinear terms may be broken down into

$$\begin{aligned} uu_x &= \sum_{n=0}^{\infty} \mathcal{A}_n, \\ vv_x &= \sum_{n=0}^{\infty} \mathcal{B}_n, \\ (uv)_x &= \sum_{n=0}^{\infty} \mathcal{C}_n, \end{aligned} \quad 41$$

where

$$\begin{aligned} \mathcal{A}_0 &= 2u_0 \frac{\partial u_0}{\partial x}, \\ \mathcal{A}_1 &= 2u_0 \frac{\partial u_1}{\partial x} + 2u_1 \frac{\partial u_0}{\partial x} \\ \mathcal{B}_0 &= 2v_0 \frac{\partial v_0}{\partial x}, \\ \mathcal{B}_1 &= 2v_0 \frac{\partial v_1}{\partial x} + 2v_1 \frac{\partial v_0}{\partial x} \\ \mathcal{C}_0 &= \frac{\partial}{\partial x} (u_0 v_0), \\ \mathcal{C}_1 &= \frac{\partial}{\partial x} (u_0 v_1 + u_1 v_0). \end{aligned} \quad 42$$

Substituting Eq.41 and Eq.40 in Eq.39,

$$\begin{aligned} \sum_{n=0}^{\infty} u_n &= \sin(x) \\ &+ E^{-1} \left\{ (1-a+aw) E \left\{ \begin{aligned} &\frac{\partial^2}{\partial x^2} (\sum_{n=0}^{\infty} u_n) \\ &+ \sum_{n=0}^{\infty} \mathcal{A}_n \\ &- \sum_{n=0}^{\infty} \mathcal{C}_n \end{aligned} \right\} \right\}, \\ \sum_{n=0}^{\infty} v_n &= \sin(x) \\ &+ E^{-1} \left\{ (1-b+bw) E \left\{ \begin{aligned} &\frac{\partial^2}{\partial x^2} (\sum_{n=0}^{\infty} v_n) \\ &+ \sum_{n=0}^{\infty} \mathcal{B}_n \\ &- \sum_{n=0}^{\infty} \mathcal{C}_n \end{aligned} \right\} \right\}. \end{aligned} \quad 43$$

When both sides of the Eq.43 are compared,

$$\begin{aligned} u_0 &= \sin(x), \\ v_0 &= \sin(x), \end{aligned}$$

$$\begin{aligned} u_1 &= E^{-1} \left\{ (1-a+aw) E \left\{ \begin{aligned} &\frac{\partial^2 u_0}{\partial x^2} \\ &+ \mathcal{A}_0 \\ &- \mathcal{C}_0 \end{aligned} \right\} \right\}, \\ &= -\sin(x) (1-a+at), \end{aligned}$$

$$\begin{aligned} v_1 &= E^{-1} \left\{ (1-b+bw) E \left\{ \begin{aligned} &\frac{\partial^2 v_0}{\partial x^2} \\ &+ \mathcal{B}_0 \\ &- \mathcal{C}_0 \end{aligned} \right\} \right\}, \\ &= -\sin(x) (1-b+bt), \end{aligned}$$

$$\begin{aligned} u_2 &= E^{-1} \left\{ (1-a+aw) E \left\{ \begin{aligned} &\frac{\partial^2 u_0}{\partial x^2} + \mathcal{A}_1 - \mathcal{C}_1 \end{aligned} \right\} \right\}, \\ &= \sin(x) [(1-a)^2 + (a-2a^2)t + \frac{1}{2}a^2t^2 \\ &\quad + 2(a-\beta) \cos x [(1-a) + (2a-1)t \\ &\quad - \frac{1}{2}a^2t^2]], \end{aligned}$$

$$v_2 = E^{-1} \left\{ (1-b+bw) E \left\{ \begin{aligned} &\frac{\partial^2 v_0}{\partial x^2} \\ &+ \mathcal{B}_1 \\ &- \mathcal{C}_1 \end{aligned} \right\} \right\},$$

$$\begin{aligned} &= \sin(x) [(1-b)^2 + (2b-2b^2)t + \frac{1}{2}b^2t^2 \\ &\quad + 2(b-a) \cos(x) [(1-b) + (2b-1)t \\ &\quad - \frac{1}{2}b^2t^2]]. \end{aligned}$$

Therefore, the approximate solution of Eq.36 is given by

$$\begin{aligned} u &= \sin(x) [(1-a+a^2) + (a-2a^2)t + \frac{1}{2}a^2t^2 \\ &\quad + 2(a-b) \cos(x) [(1-a) + (2a-1)t \\ &\quad - \frac{1}{2}a^2t^2] + \dots], \\ v &= \sin(x) [(1-b+b^2) + (b-2b^2)t \\ &\quad + \frac{1}{2}b^2t^2 + 2(b-a) \cos(x) [(1-b) + (2b- \\ &\quad 1)t - \frac{1}{2}b^2t^2] + \dots]. \end{aligned} \quad 44$$

If put $a \rightarrow 1$ and $b \rightarrow 1$ in Eq.44, the problem solution will be re-created as follows:

$$\begin{aligned} u(x, t) &= \sin(x) \left(1 - t + \frac{t^2}{2!} - \dots \right), \\ v(x, t) &= \sin(x) \left(1 - t + \frac{t^2}{2!} - \dots \right). \end{aligned} \quad 45$$

This is the closed form equivalent of the precise solution:

$$\begin{aligned} u(x, t) &= \sin(x) e^{-t}, \\ v(x, t) &= \sin(x) e^{-t}. \end{aligned} \quad 46$$

Table 3 displays the values of the exact and approximate solutions of system 36 at the different a and b values, also includes the absolute error of the approximate solutions for the exact solutions. Fig. 3 shows graphs of the exact and approximate solutions of system 36 at the different a and b values between 0 and 1.

Table 3. The approximate and exact solution of system 36 with fractional operator CF, u_1 is app. sol. of u, v at $a = 0.9$, u_2 is app. sol. of u, v at $a = 1$, u_1 is exact sol. of u, v .

x	t	u_1	u_2	u_3	$ u_1 - u_3 $	$ u_2 - u_3 $
0.2500	0.2500	0.1869	0.1933	0.1927	0.0058	0.0006
0.5000	0.2500	0.3621	0.3746	0.3734	0.0113	0.0012
0.7500	0.2500	0.5149	0.5325	0.5309	0.0160	0.0017
0.2500	0.5000	0.1611	0.1546	0.1501	0.0111	0.0046
0.5000	0.5000	0.3122	0.2996	0.2908	0.0214	0.0089
0.7500	0.5000	0.4439	0.4260	0.4134	0.0305	0.0126
0.2500	0.7500	0.1479	0.1314	0.1169	0.0310	0.0146
0.5000	0.7500	0.2866	0.2547	0.2265	0.0601	0.0282
0.7500	0.7500	0.4075	0.3621	0.3220	0.0855	0.0401

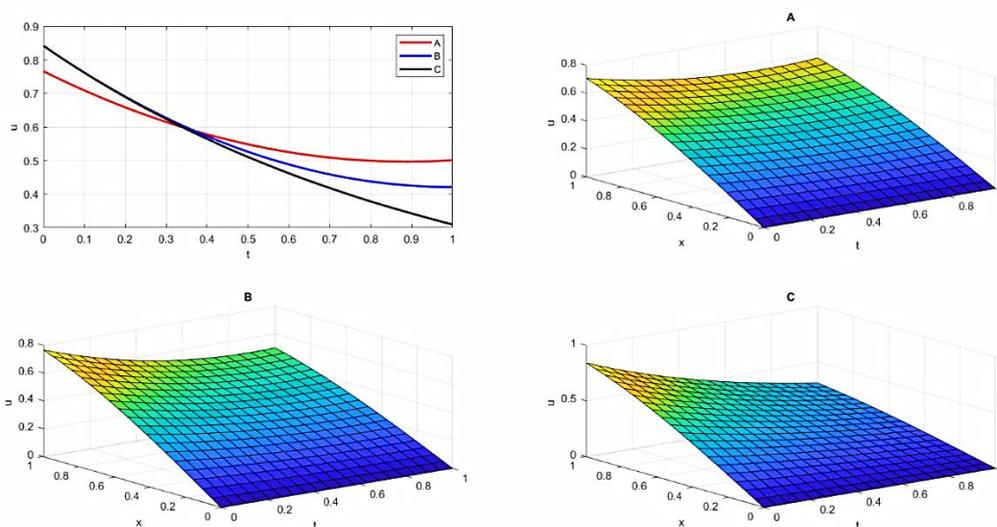


Figure 3. The graphs of the approximate and the exact solutions of nonlinear system 36 among different values of x and t in case (3D) and x is fixed in case (2D) when $a = 0.9, 1$ for nonlinear system in the CF fractional operator, (A) is the app. sol. of u, v at $a = 0.9$, (B) is the app. sol. of u, v at $a = 1$, (C) is the exact sol. of u, v .

Remark. By comparing the results of this method with Yang decomposition method it appears that the results are similar to both methods²⁰.

Conclusion

In this article, the Elzaki decomposition method has been presented in terms of its derivation and convergence and its application to fractional differential equations (FDEs). The method was convergent and efficient to solve fractional differential equations with the Caputo-Fabrizio

operator. The approximate solution of fractional differential equations (FDEs) with the derivative of Caputo-Fabrizio was convergent to the exact solution. Finally, this method can be adopted to solve fractional differential equations of this type.

Acknowledgment

We appreciate the efforts of everyone who contributed even a little to this work.

Authors' Declaration

- Conflicts of Interest: None.
- I hereby confirm that all the Figures and Tables in the manuscript are mine. Furthermore, any Figures and images, that are not mine, have been

- included with the necessary permission for republication, which is attached to the manuscript.
- Ethical Clearance: The project was approved by the local ethical committee in University of Thi-Qar.

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استخدام طريقة تحليل الزاكي لحل المعادلات التفاضلية غير الخطية ذات الرتبة الكسرية مع المؤثر الكسري كابوتو-فابريزيو

محمد عبد الشريف حسين

مركز البحث العلمي، جامعة العين، ذي قار، العراق

الخلاصة

يتم تطبيق تقنيات حساب التفاضل والتكامل الجزئي بنجاح في العديد من فروع العلوم والهندسة ، احدى هذه التقنيات طريقة تحليل الزكي وادوميان (EADM) ، والتي لم يدرسها الباحثون باستخدام المشتق الكسري لـ Caputo-Fabrizio. يهدف هذا العمل لدراسة طريقة تحليل الزكي وادوميان (EADM) لحل المعادلات التفاضلية الكسرية بمشتق Caputo-Fabrizio. قدمنا خوارزمية هذه الطريقة مع المؤثر CF وناقشنا تقاربها باستخدام طريقة سلسلة كوشي بعد ذلك ، تم تطبيق الطريقة لحل معادلات برجر ، الشبيهة بالحرارة ، و برجر المزدوجة مع المؤثر Caputo -Fabrizio. في الختام أن الطريقة كانت متقاربة وفعالة لحل هذا النوع من المعادلات التفاضلية الكسرية.

الكلمات المفتاحية: معادلة برجر ، المؤثر الكسري كابوتو-فابريزيو ، طريقة تحليل الزكي ، المعادلة الشبيهة بالحرارة.